

Boundary control in transport and diffusion equations

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High Contrast - High Resolution Inverse problems

High *contrast* comes from optical/elastic parameters.

High *resolution* comes from MRI/Ultrasound.

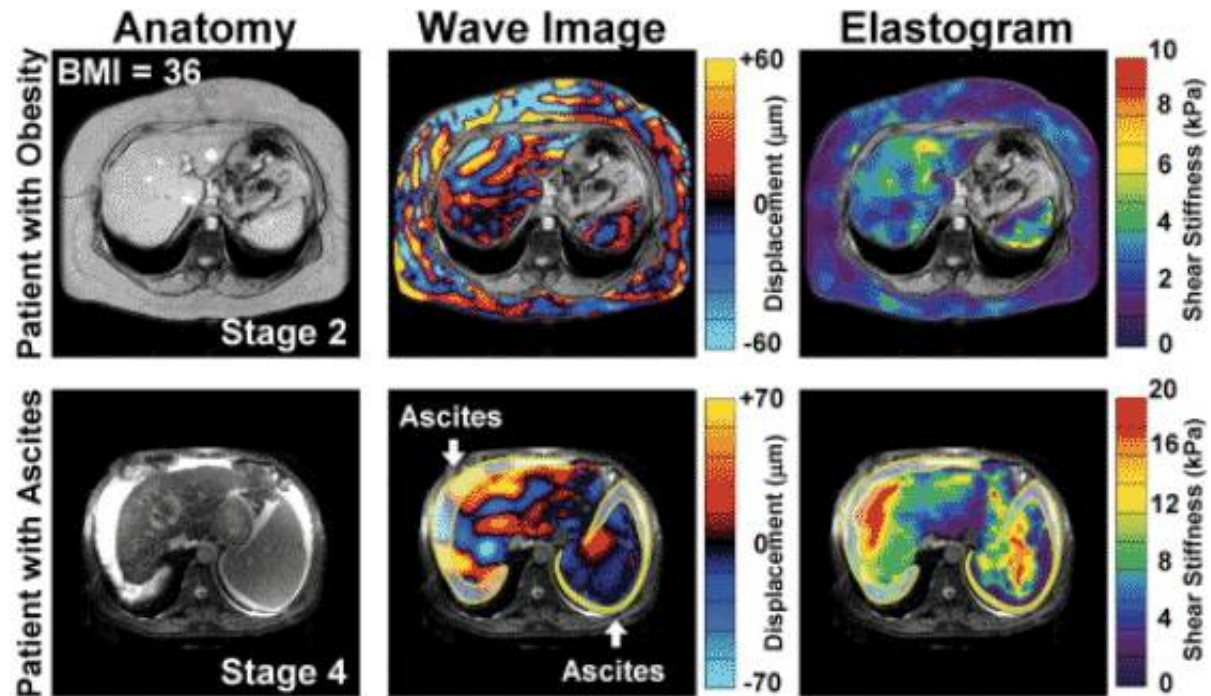
Assume *physical coupling* between high contrast and high resolution.

Inversion of the *high resolution* modality (MRI/Ultrasound) provides *internal information* about optical/elastic parameters.

Reconstruction of optical/elastic parameters from *internal information* requires diffusion solutions to satisfy **qualitative properties**.

Boundary controls are sought to ensure **qualitative properties** hold.

Magnetic Resonance Elastography

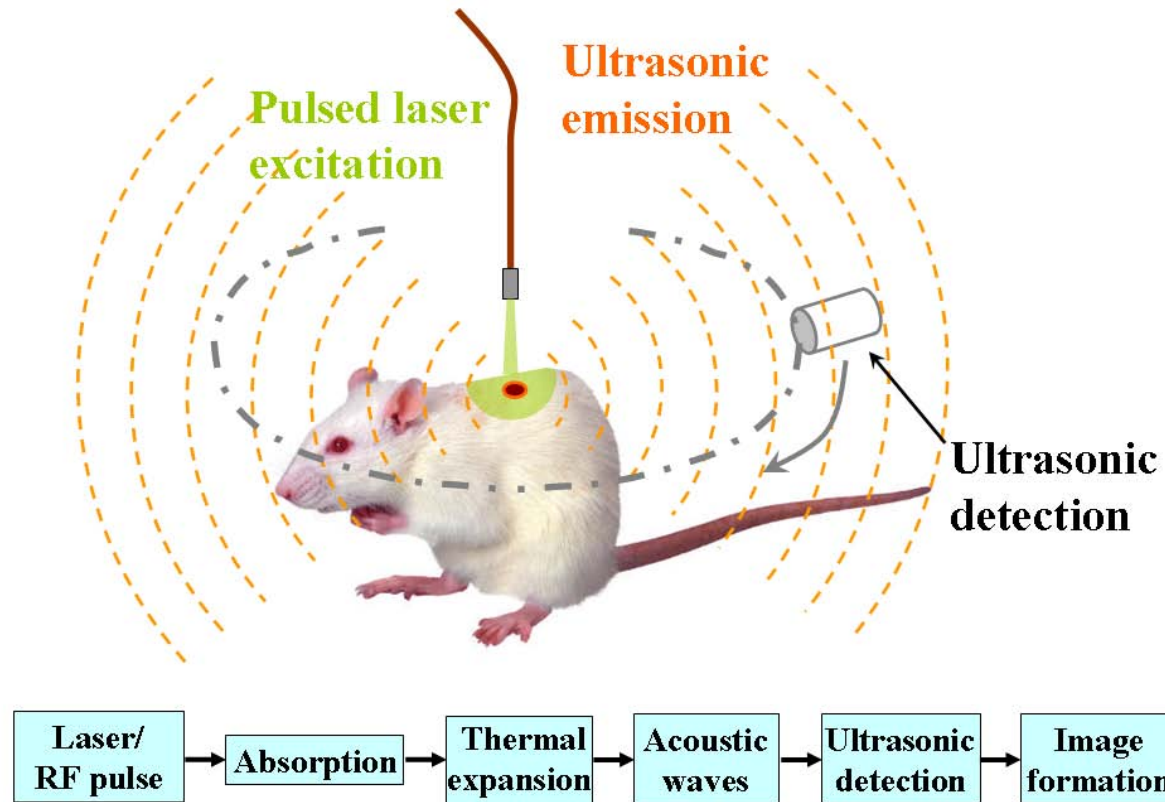


Hepatic Fibrosis by Liver Stiffness (Mayo Clinic)

Coupling between **Elastic Waves** and **Magnetic Resonance Imaging**.

MRI provides $u =$ **Elastic Displacement**.

The Photo-acoustics Effect



Coupling between Optical Radiation and Ultrasound

Inverse Ultrasound provides $\sigma u = \text{absorption} \times \text{energy density}$.

Quantitative reconstruction

The above and many other examples may be recast as the **reconstruction of coefficients in a second-order elliptic equation from knowledge of some of its solutions** $\{u_j\}_{1 \leq j \leq J}$:

$$a_{ij} \partial_{ij}^2 u_k + b_i \partial_i u_k + c u_k = 0, \quad x \in X; \quad u_k = g_k, \quad x \in \partial X.$$

Each u_k provides a *linear constraint* on (a, b, c) .

Can we find $\{g_k\}_{1 \leq k \leq J}$ such that $\{(\partial_{ij}^2 u_k, \partial_i u_k, u_k)\}_{1 \leq k \leq J}$ are of **maximal rank** compatible with above constraint in $Y \subset X$?

Locally, maximal rank family may be constructed by plane waves or generalizations of harmonic polynomials. Can such local solutions be *controlled* from *the boundary* ∂X ?

Runge Property and Boundary Control

Approximate boundary control amounts to verifying the **Runge approximation property**: Let $X_1 \subset X_2$ two (simply connected) open domains. Solutions in X_2 restricted to X_1 are *dense* in the set of solutions in X_1 . Traces of solutions on the larger domain on ∂X_2 provide appropriate boundary controls for solutions in X_1 .

The **Runge approximation** is *equivalent* for second-order elliptic operators to the **Unique Continuation Property** (UCP): two solutions of an elliptic equation on X_2 and equal on an open $X_1 \subset X_2$ have to be equal on X_2 . [Lax CPAM 56]

In the context of hybrid inverse problems for scalar equations and systems of second-order equations, see [B. InsideOut 2013]; [B. Uhlmann CPAM 2013]; [B. Monard Uhlmann SIAP 2016].

Local, global, uniform controls

The *Runge Approximation Property* allows us to **control a property of interest** (such as, e.g., linear independence of gradients of solutions) **locally**.

Can we get **global** control (over the whole domain X) ? In general, we do not know how to do so, unless explicit global solutions (harmonic polynomials, CGO solutions) can be constructed.

The control depends on the geometry (coefficients a_{ij} , domain X). Can we get a control (a choice of **boundary conditions**) **uniform** in the **coefficients**? **Not in general**: for instance, in dimension $n \geq 3$, we can construct for any given g on ∂X a (scalar) conductivity $a = a(x)$ such that $\nabla u(x_0) = 0$ for some x_0 in the vicinity of a chosen point [Alberti B. Di Cristo, 2017; Capdeboscq 2016].

[B. InsideOut 2013]

Hybrid inverse problem and transport equation

Assume that (optical) radiation is now modeled by a **transport equation** posed on a convex bounded domain $X \subset \mathbb{R}^n$; $V = \mathbb{S}^{n-1}$:

$$v \cdot \nabla u + \sigma(x, v)u = \int_V k(x, v', v)u(x, v')dv', \quad (x, v) \in X \times V$$

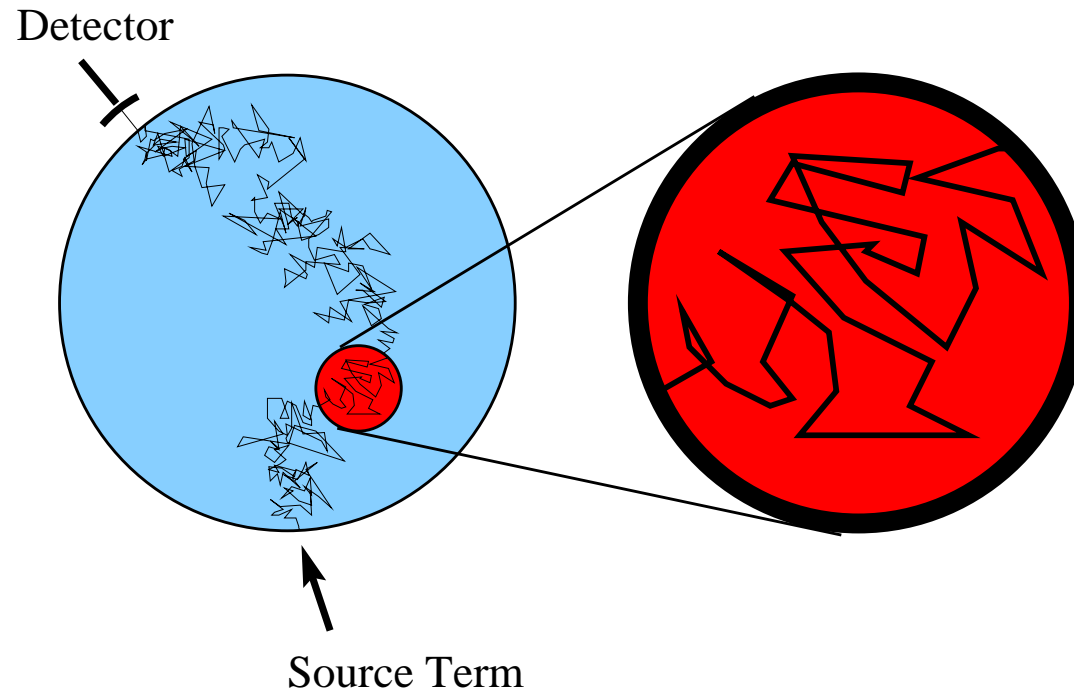
$$u = g, \quad (x, v) \in \Gamma_-; \quad \Gamma_{\pm} = \{(x, v) \in \partial X \times V, \pm v \cdot n(x) > 0\}.$$

Assume that optical radiation is coupled to a **high resolution** modality (ultrasound) that **provides** $u(x, v)$. [B. Chung Schotland 2015]

Optical coefficients $(\sigma(x, v), k(x, v', v))$ are then fully reconstructed provided $u(x, v)$ can locally (in x) take any **prescribed v -dependent shape**.

This raises the question of **boundary control of transport solutions**.

Transport versus Diffusion



The solution of the **transport equation** $v \cdot \nabla u_\varepsilon + \frac{1}{\varepsilon} \sigma_s(x)(u_\varepsilon - \int_V u_\varepsilon dv) + \varepsilon \sigma_a(x)u_\varepsilon = 0$ converges as $\varepsilon \rightarrow 0$ to the solution of the **diffusion equation** (elliptic second-order equation) : $-\nabla \cdot \frac{1}{n\sigma_s(x)} \nabla U + \sigma_a(x)U = 0$.

Control problems

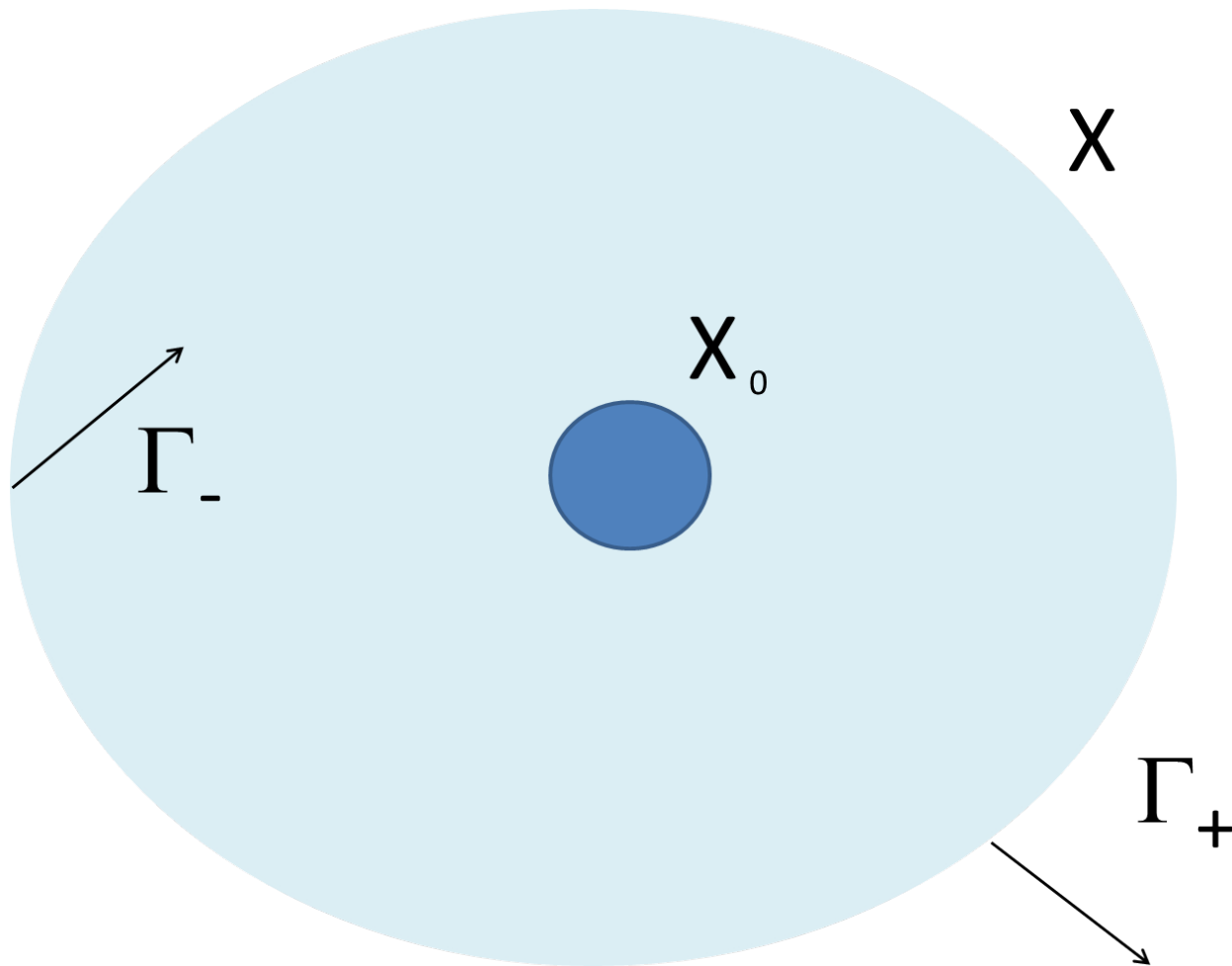
Coupled physics inverse problems require **boundary controls**.

For **diffusion equations**, this is possible thanks to the **unique continuation principle**.

Transport solutions in the high scattering (small mean free path) regime are well approximated by **diffusion solutions**.

- Can we **control** transport solutions from the boundary?
- Is this related to a **unique continuation principle**?
- Is the latter even valid?

Geometry of transport problem



Control of Outgoing Transport Solution

$$Tu := v \cdot \nabla u + \sigma(x, v)u = \int_V k(x, v', v)u(x, v')dv' =: Ku, \quad (x, v) \in X \times V$$

$$u = g, \quad (x, v) \in \Gamma_-; \quad \Gamma_{\pm} = \{(x, v) \in \partial X \times V, \pm v \cdot n(x) > 0\}.$$

Assume that $0 \leq \int_V k(x, v', v)dv' \leq \sigma(x, v)$ (**sub-criticality** condition).

Then **albedo** $\mathcal{A} : g \mapsto u|_{\Gamma_+}$ **well posed** from $L^2(\Gamma_-, d\xi)$ to $L^2(\Gamma_+, d\xi)$.

Control of outgoing solutions:

Can we find g on Γ_- such that $u|_{\Gamma_+}$ is prescribed?

More precisely, what is the Range of \mathcal{A} ?

Fredholm theory

Using a *contraction argument* or a *compact perturbation argument*, we obtain that \mathcal{A} is a **Fredholm operator** with vanishing index so that:

$$L^2(\Gamma_+, d\xi) = \text{Ran}(\mathcal{A}) \oplus \text{Ker}(\mathcal{A}^*), \quad L^2(\Gamma_-, d\xi) = \text{Ran}(\mathcal{A}^*) \oplus \text{Ker}(\mathcal{A}).$$

Any g with a non-trivial component in $\text{Ker}(\mathcal{A}^*)$ therefore cannot be *controlled* from Γ_- .

Note that any control g on Γ_- in $\text{Ker}(\mathcal{A})$ leaves no trace on Γ_+ !

Integral formulation and contraction argument

Consider the problem

$$v \cdot \nabla u + \sigma u =: Tu = Ku, \quad (x, v) \in X \times V, \quad u = g, \quad (x, v) \in \mathcal{C}.$$

Here \mathcal{C} is either Γ_{\pm} or any conditions such that $Tu = 0$ with prescribed (Dirichlet) conditions on \mathcal{C} is invertible. Then the above is equivalent to

$$u = T_{\mathcal{C}}^{-1}Ku + L_{\mathcal{C}}g.$$

For k and hence K sufficiently small, the above problem with $\mathcal{C} = \Gamma_{+}$ admits a unique solution and $\mathcal{A} : L^2(\Gamma_{-}, d\xi) \rightarrow L^2(\Gamma_{+}, d\xi)$ is invertible.

For sufficiently small scattering k , the boundary control is *exact* with $L^2(\Gamma_{+}, d\xi) = \text{Ran}(\mathcal{A})$.

Compactness argument

Under mild assumptions on K , we show following [Mikhlin; Mokhtar-Kharroubi; GLPS 88] that $T_{\mathcal{C}}^{-1}K$ is a compact operator. Using T_{\pm}^{-1} for $\mathcal{C} = \Gamma_{\pm}$, let

$$N = \text{Ker}(I - T_{+}^{-1}K) \neq \{0\}, \quad N_{-} = \{u|_{\Gamma_{-}}, u \in N\} = \text{Ker}(\mathcal{A}).$$

Note $\text{Ker}(I - T_{-}^{-1}K) = \{0\}$ by sub-criticality and hence $\dim N = \dim N_{-}$.

Then identify compatibility conditions to solve $(I - T_{+}^{-1}K)v = L_{+}g$ with those to solve $\mathcal{A}^*h = g$ to obtain that $L^2(\Gamma_{-}, d\xi) = \text{Ran}(\mathcal{A}^*) \oplus \text{Ker}(\mathcal{A})$. Pass to adjoints for decomposition of $L^2(\Gamma_{+}, d\xi)$.

It remains to understand whether $\mathcal{N} = \dim N > 0$ is possible.

Non-controllability

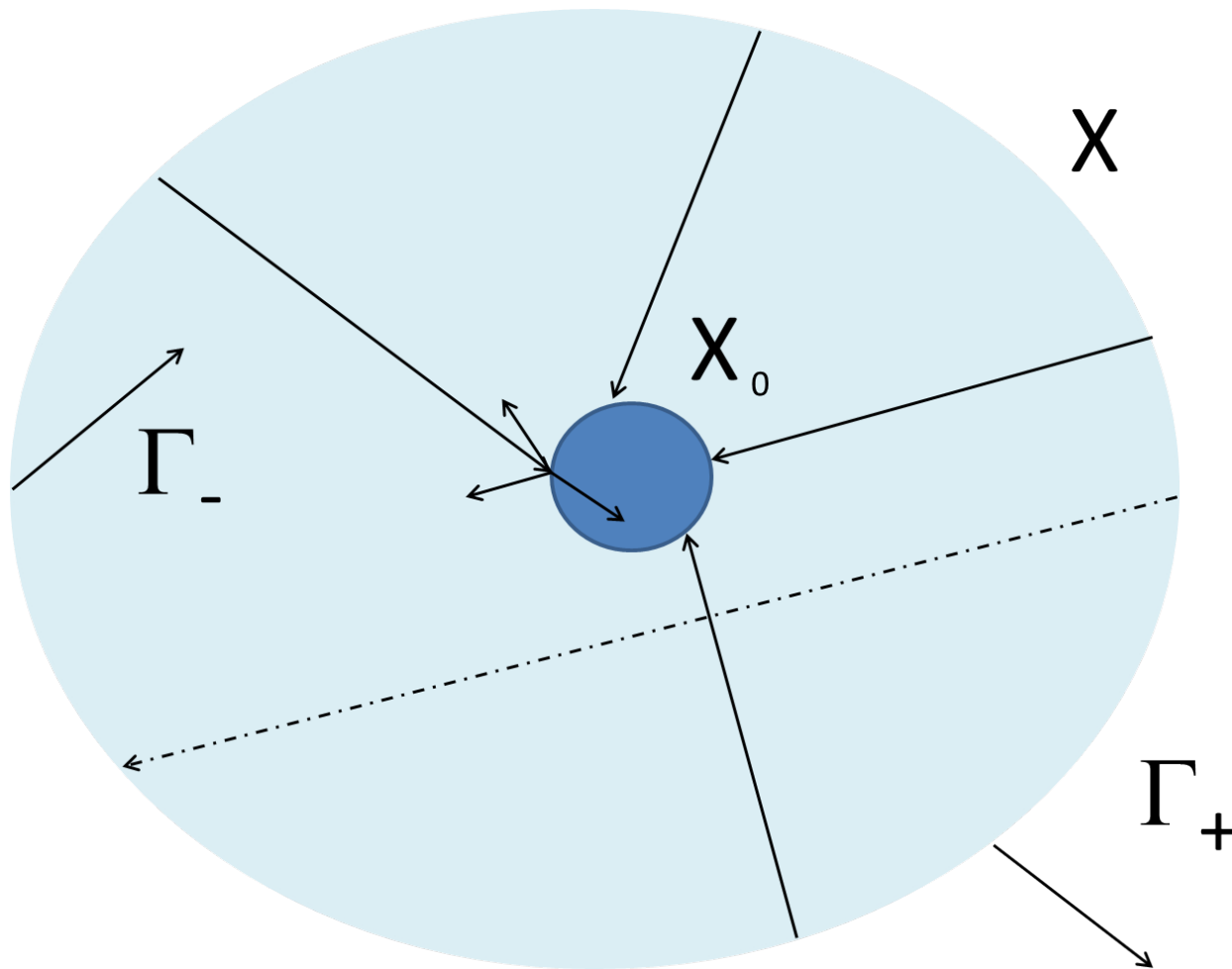
There are coefficients $k(x) \leq \sigma(x)$ such that $\mathcal{N} = \dim \text{Ker}(I - T_+^{-1}K) > 0$.

This implies that $\dim \text{Ker}(\mathcal{A}) = \dim \text{Ker}(\mathcal{A}^*) = \mathcal{N} > 0$.

This implies the existence of $g \neq 0$ on Γ_- such that $\mathcal{A}g = 0$ on Γ_+ : **this is a non-radiating boundary condition.**

This implies the existence of \mathcal{N} **compatibility conditions** on h defined on Γ_+ such that h can be *controlled* by some g on Γ_- , i.e., $h = \mathcal{A}g$. **Not every outgoing boundary condition can be controlled.**

Geometry of internal control problem



Internal controls

Let us come back to the **internal control problem** and consider X *strictly convex* and $X_0 \Subset X$ also *strictly convex*. Let u_0 be a solution of

$$v \cdot \nabla u_0 + \sigma u_0 = \int_V k(x, v', v) u_0(x, v') dv' \quad X_0 \times V$$

and let g_0 denote its trace on $\Gamma_+(X_0) \cup \Gamma_-(X_0)$. (essentially $\partial X_0 \times V$.)

Can we find u solution to the following problem?

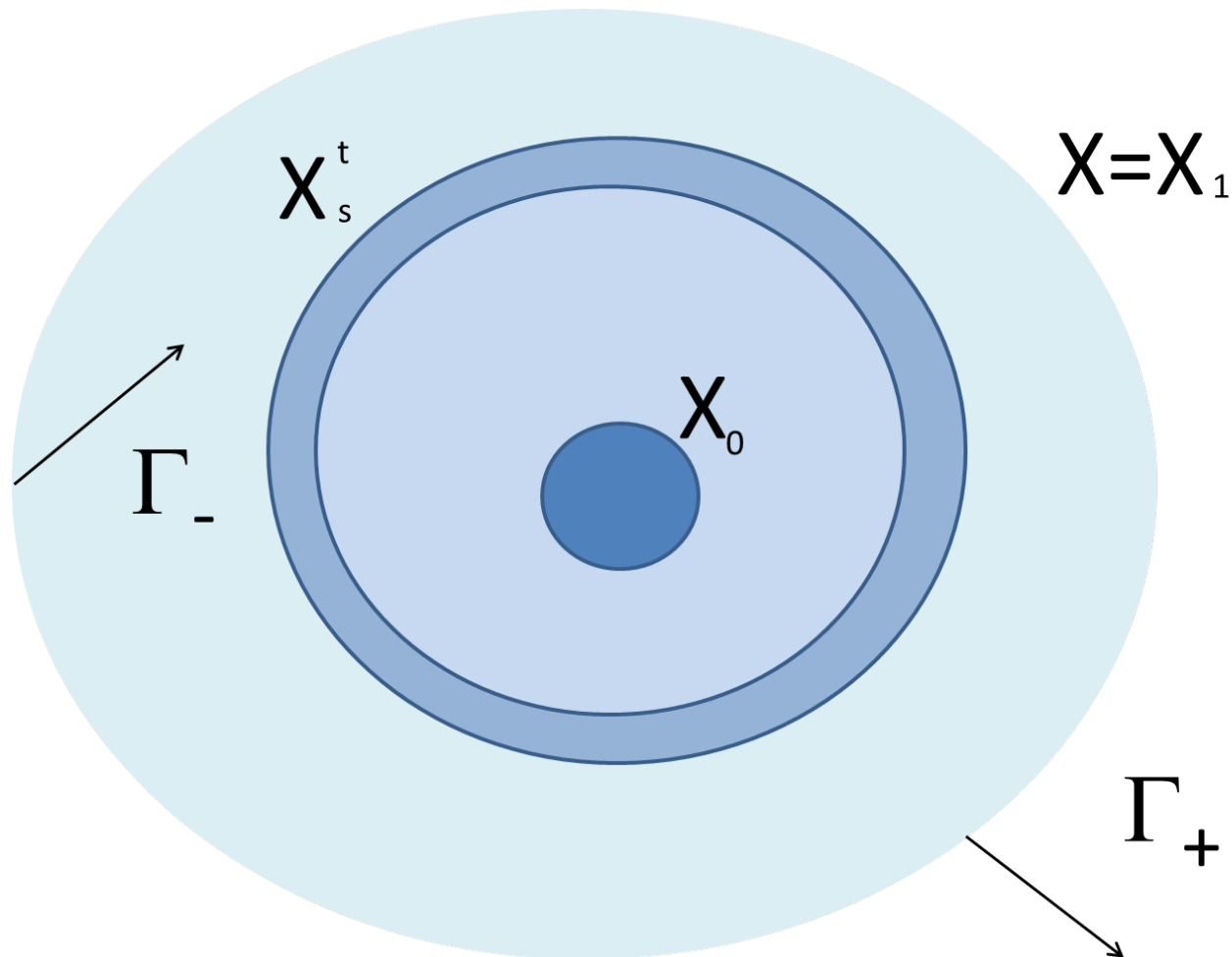
$$\begin{aligned} v \cdot \nabla u + \sigma u &= \int_V k(x, v', v) u(x, v') dv' \quad X_0^1 \times V, \quad X_0^1 = X \setminus \overline{X_0} \\ u &= g_0 \text{ on } \Gamma_+(X_0) \cup \Gamma_-(X_0). \end{aligned}$$

Theorem. The answer is positive. The solution is not unique.

Joint with Alexandre Jollivet; see also [B. Chung Schotland 2015].

For $g = u|_{\Gamma_-(X)}$, we obtain that $u|_{X_0}$ is equal to the prescribed u_0 . We have an **exact internal control**.

Derivation by layer peeling argument



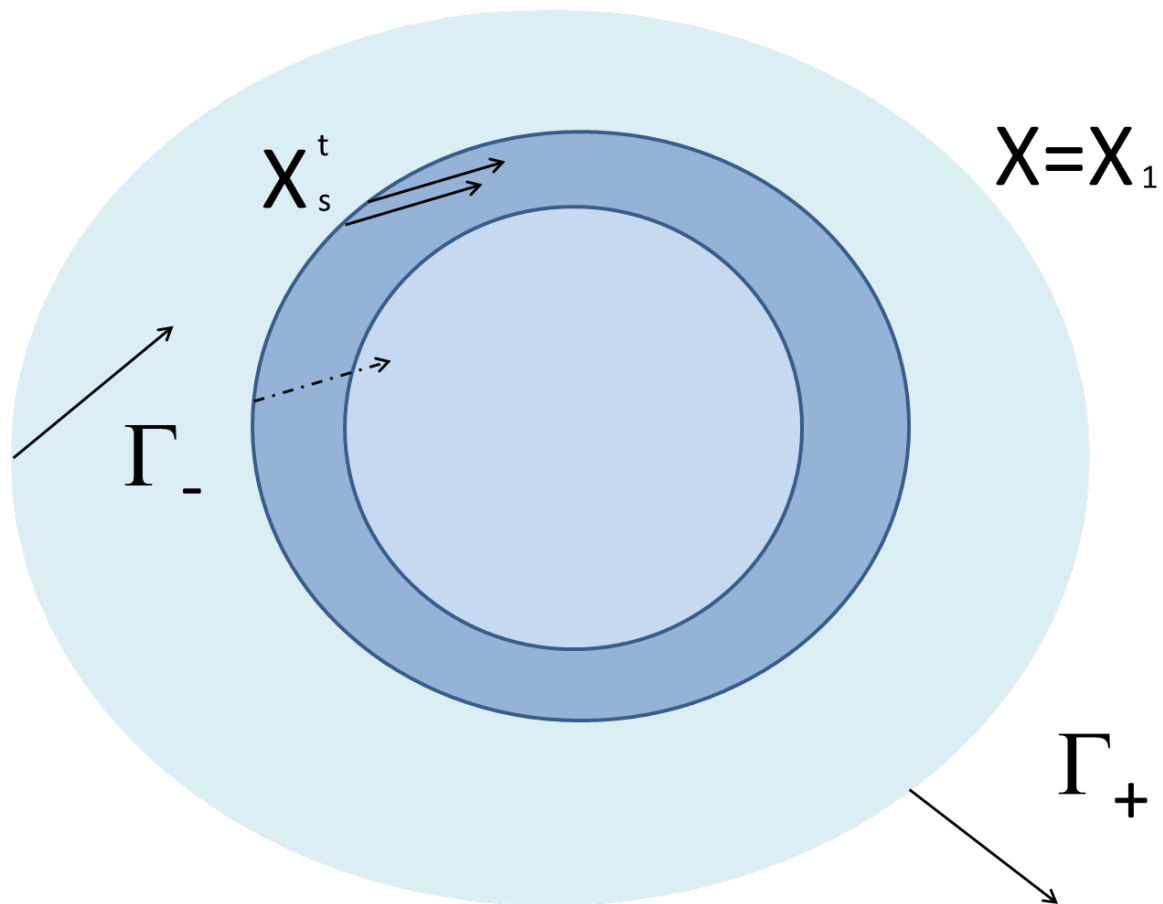
No Unique continuation in transport

Assume u solution on X with $u = 0$ on X_0 . Then is $u = 0$ on X ?

Result. The answer is negative: **UCP does not hold for transport.**

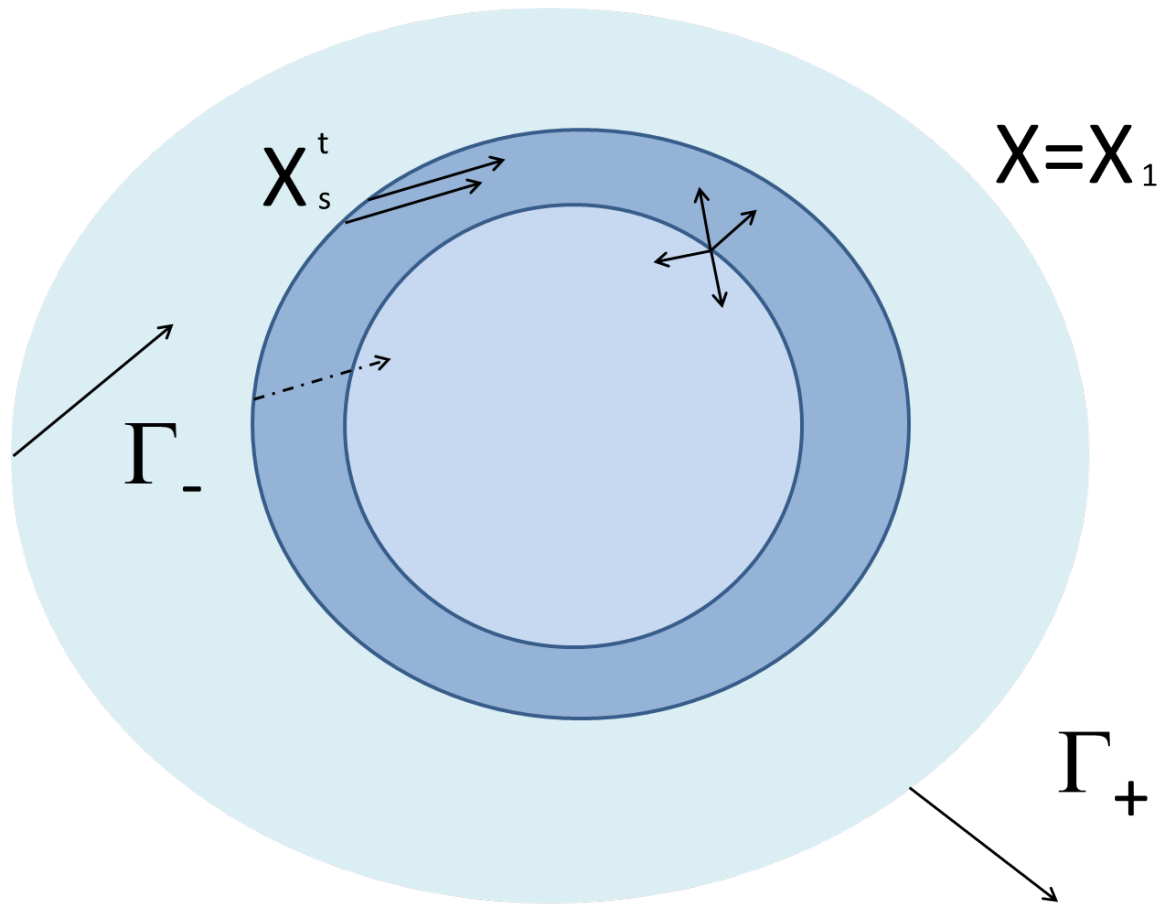
Consider the last layer $X_{s_{J-1}}^{s_J}$ and consider $h \neq 0$ with support in $\Gamma_-(X)$ on the set of incoming conditions that do not meet $Y_{s_{J-1}}$. Let u_h be the solution on X_0^1 with such a boundary condition and vanishing incoming conditions on Y_0 . Let g_0 be the trace of u_h on ∂X_0 . Let u be constructed as in the preceding result so that in particular $u = 0$ on the support of h . Consider $v = u - u_h$. Then $v = 0$ on ∂X_0 and is extended by 0 on X_0 . v solves the transport equation on X and $v = -h$ on the support of h in $\Gamma_-(X)$. As a consequence, the transport solution v cannot vanish on the whole of X even though it vanishes on $X_0 \times V$.

No Unique Continuation Property



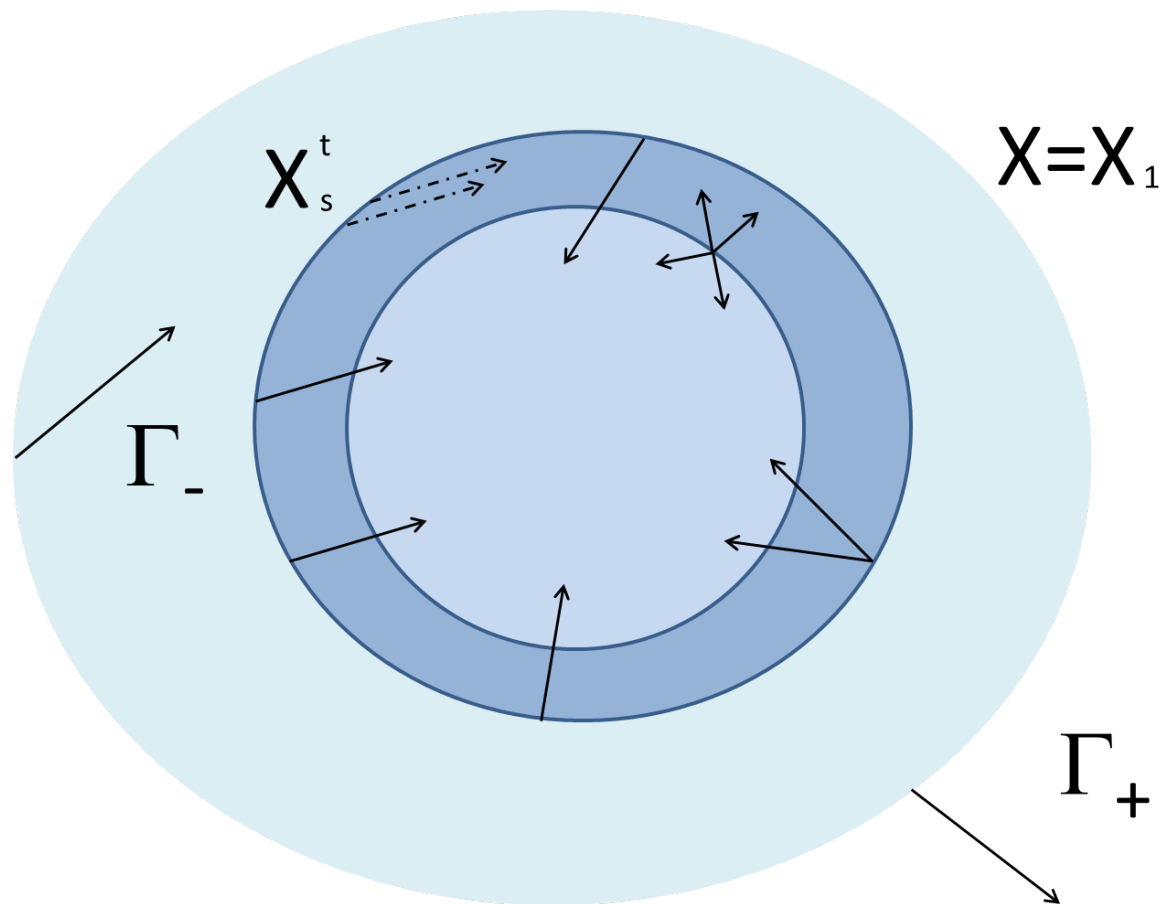
Consider an incoming condition on $\Gamma_-(X_t)$ on lines not crossing X_s

No Unique Continuation Property



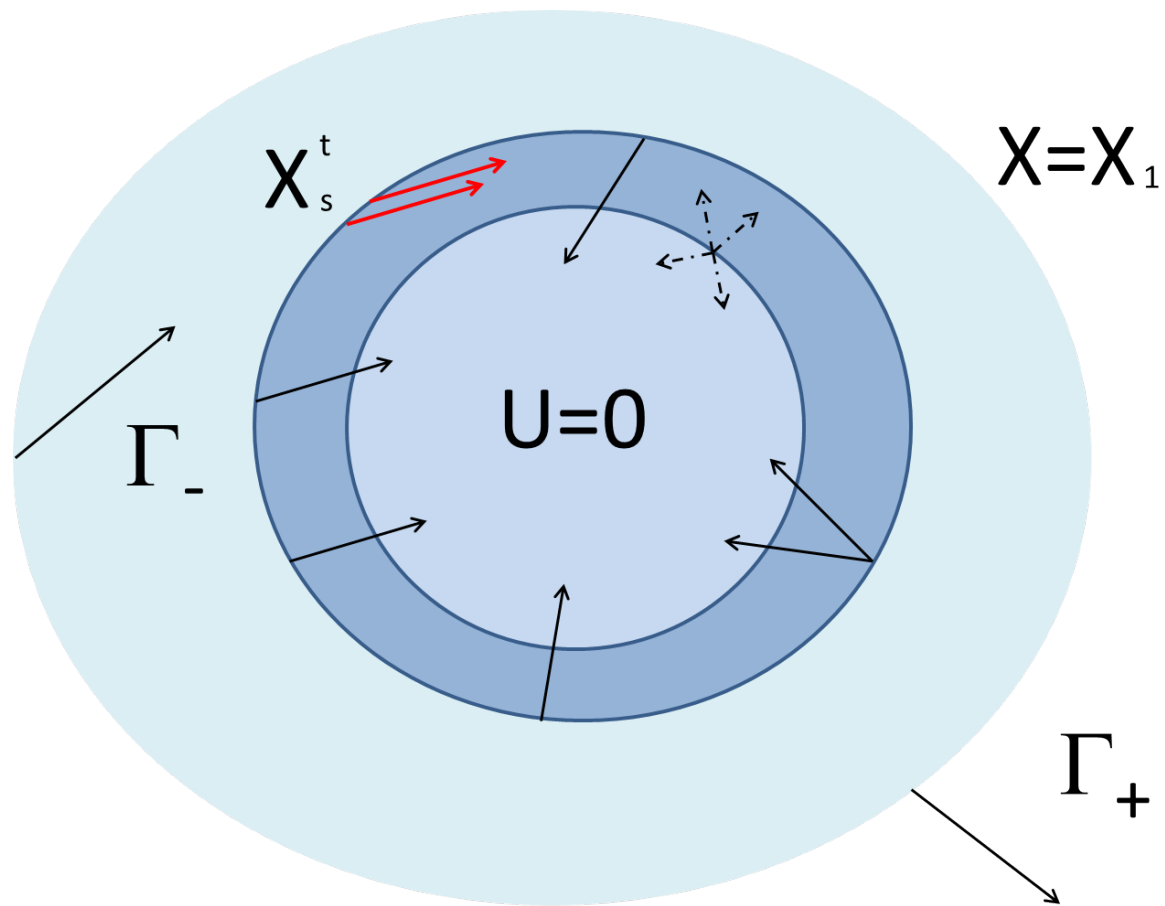
Solve Transport on large domain X_t and identify trace on $\partial X_s \times V$.

No Unique Continuation Property



Control trace on $\partial X_s \times V$ from boundary ∂X_t : lines not crossing X_s not necessary.

No Unique Continuation Property



Add negative of initial radiation condition on ∂X_t to violate UCP.

Conclusions

Heuristically, in an inverse problems/imaging method, we wish the *probing signal* to be as "varied" and "informative" as possible.

Realizing this variability from the available part of the boundary amounts to a **boundary control** problem.

For diffusion equations, **approximate** boundary control is possible as an application of a **unique continuation property**. Such a control is an *unstable* process.

For transport equations, **exact** boundary control is feasible even though **UCP does not hold**. However, the stability degrades as the mean free path tends to 0.