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## Euler's eqs for body, ang mom of a free rigid body

Given a symm, pos def  $3 \times 3$  matrix  $\mathbb{I}$ ,  
def. the 1<sup>st</sup> order ODE

$$\dot{\underline{M}} = \underline{M} \times (\mathbb{I}^{-1} \underline{M}) \quad \text{on } \mathbb{R}^3$$

Rel to rigid body? Later.

Def  $E: \mathbb{R}^3 \rightarrow \mathbb{R}$  by

$$E(\underline{M}) := \frac{1}{2} \langle \underline{M}, \mathbb{I}^{-1} \underline{M} \rangle$$

Properties: (a) If  $\underline{M}(t)$  is a sol curve,

(i)  $\|\underline{M}(t)\| = \text{const}$

(ii)  $E(\underline{M}(t)) = \text{const}$

(iii)  $\dot{\underline{M}} = \underline{M} \times \nabla E(\underline{M})$

Pf: (i) Obvious

(ii) Prove for  $\dot{\underline{M}} = \underline{M} \times \nabla E(\underline{M})$ ,  $E$  is a  
diff fun on  $\mathbb{R}^3$  (Ham wrt Poisson structure on  $\mathbb{R}^3$ )

$$\begin{aligned} \frac{d}{dt} \|\underline{M}(t)\|^2 &= 2 \langle \underline{M}(t), \dot{\underline{M}}(t) \rangle \\ &= \langle \underline{M}(t), \underline{M} \times \nabla E(\underline{M}) \rangle = 0 \end{aligned}$$

$$\begin{aligned} \frac{d}{dt} E(\underline{M}(t)) &= \langle \nabla E(\underline{M}(t)), \dot{\underline{M}}(t) \rangle \\ &= \langle \nabla E, \underline{M} \times \nabla E \rangle = 0. \end{aligned}$$

Consequences? Trajs lie in intersections of spheres w/ level sets of  $E$ .

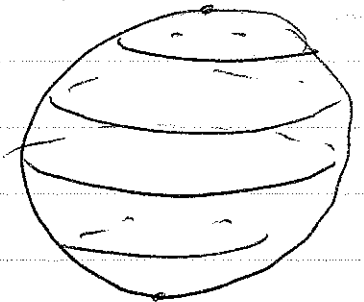
Case 1:  $\mathbb{I} = \lambda \mathbb{I}$  for some  $\lambda \in \mathbb{R}^+$

$$\dot{M} = M \times \left(\frac{1}{\lambda} M\right) = Q$$

$\Rightarrow$  all pts are eq.

Case 2:  $\mathbb{I} = \text{diag}(\lambda_1, \lambda_1, \lambda_3)$ ,  $\lambda_1 \neq \lambda_3$   
mod change of basis.

Level sets of  $E$  are 'upright' spheroids, passing through  $Q$ , w/ vertical axis of symm. Intersections w/ spheres are horiz circles, plus poles.



Equator + poles are eq;  
all pts in N/S hemisphere rotate in same direction

$$\text{Explicit calc: } \mathbb{I} = \lambda_1 \mathbb{I} + (\lambda_3 - \lambda_1) e_3 e_3^T$$

$$\mathbb{I}^{-1} = \frac{1}{\lambda_1} \mathbb{I} + \left(\frac{1}{\lambda_3} - \frac{1}{\lambda_1}\right) e_3 e_3^T$$

$$\Rightarrow \dot{M} = \underbrace{\left(\frac{1}{\lambda_1} - \frac{1}{\lambda_3}\right)}_{\text{sign det's dir of rotation}} \langle M, e_3 \rangle e_3 \times M$$

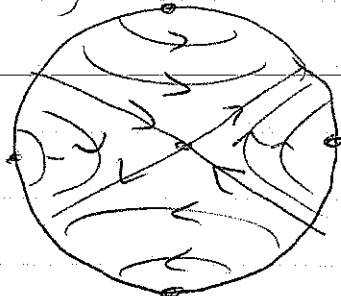
sign det's dir of rotation  
in each hemisphere.

$M$  rotates w/ const speed, since  $\langle M, e_3 \rangle$  const along traj.

Case 3:  $\mathbb{I} = \text{diag}(\lambda_1, \lambda_2, \lambda_3)$ ,  $\lambda_1 > \lambda_2 > \lambda_3$   
mod change of basis.

Level sets of  $E$  are triaxial ellipsoids.

Fixing a sphere & varying ellipsoids, find six pts of intersection, remaining intersections are curves:



Pts are eq - eigenvalues of  $\mathbb{I}$ .

~~Can show 2 saddles (vectors of  $\lambda_2$ ), four centers: lin stab?~~

$$\text{let } X(M) := \underline{M} \times (\mathbb{I}^{-1} \underline{M})$$

$$\text{if } \mathbb{I} \underline{M}_e = \lambda_j \underline{M}_e, \text{ then } X(\underline{M}_e) = 0.$$

let  $\hat{\cdot} : \mathbb{R}^3 \rightarrow$  skew-sym  $3 \times 3$  matrices:

$$\hat{y} z = y \times z \quad \forall y, z \in \mathbb{R}^3$$

$$X(M) = \hat{M} \mathbb{I}^{-1} M = -\mathbb{I}^{-1} M M^T$$

$$\nabla X(M_e) = \hat{M}_e \mathbb{I}^{-1} - \mathbb{I}^{-1} M_e$$

$$= \hat{M}_e \mathbb{I}^{-1} - \frac{1}{\lambda_j} \hat{M}_e$$

$$= \hat{M}_e (\mathbb{I}^{-1} - \frac{1}{\lambda_j} \mathbb{I})$$

$$= \hat{M}_e \text{diag} \left( \frac{1}{\hat{\alpha}_1} - \frac{1}{\hat{\alpha}_j}, \quad \text{~~1/\hat{\alpha}_2 - 1/\hat{\alpha}_j~~, \quad \frac{1}{\hat{\alpha}_2} - \frac{1}{\hat{\alpha}_j}, \quad \frac{1}{\hat{\alpha}_3} - \frac{1}{\hat{\alpha}_j} \right)$$

0 is an eigenvalue of vector  $\hat{M}_e$  - there's no motion off the sphere!

If  $j=1$  or  $3$ , get conj pair of imag eigenvalues; if  $j=2$ , get real pair summing to 0.

### Restriction of Euler eqn to sphere

We can regard  $\dot{M} = M \times \nabla E(M)$  as an ODE on a sphere, since  $\|M\| = \text{const}$ .

Advantages: i) lower dim

ii) 'boring' direction is eliminated (for  $E(M) = \frac{1}{2} \langle M, M^{-1} M \rangle$ ,  $\Pi$  w/ distinct eigenvalues, zero eigenvalues of  $\dot{M}$  at eq. are eliminated)

~~Advantages~~ iii) 2-spheres are symplectic manifolds.

Dis-adv: i) nonlinear manifold

ii) signif of  $\mathbb{R}^3 \sim \mathfrak{so}(3)$  is possibly obscured (next example)

## Full free rigid body exp

Quick review of rotation group  $SO(3)$ :

$$SO(3) = \{ A \in \mathbb{R}^{3 \times 3} : A^T A = \mathbb{1}, \det A = 1 \},$$

i.e. orth + or. - pres.

$SO(3)$  is a non-lin. manifold, +  
matrix group.

Tangent space?

Given a curve  $A(\epsilon)$  in  $SO(3)$ ,

$$\mathbb{1} \equiv A(\epsilon)^T A(\epsilon)$$

$$\Rightarrow 0 = (A'(0))^T A(0) + A(0)^T A'(0)$$

$$\text{Let } A_0 = A(0), B = A_0^{-1} A'(0) = A_0^T A'(0).$$

$$\begin{aligned} 0 &= (A_0 B)^T A_0 + A_0^T (A_0 B) \\ &= B^T (A_0^T A_0) + (A_0^T A_0) B \\ &= B^T + B \end{aligned}$$

$\Rightarrow B$  is skew-symmetric.

On the other hand, given  $B$  skew-symm.,  
def.  $A(\epsilon) := A_0 \exp(\epsilon B)$ .

$$\begin{aligned} A(\epsilon)^T A(\epsilon) &= (\exp(\epsilon B))^T A_0^T A_0 \exp(\epsilon B) \\ &= \exp(\epsilon B^T) \exp(\epsilon B) \end{aligned}$$

$$= \exp(-\epsilon B) \exp(\epsilon B)$$

$$= I$$

$A(\epsilon)$  orth  $\Rightarrow |\det A(\epsilon)| = 1$ ;  
 $\det A(0) = \det A_0 = 1$ , so  $A(\epsilon) \in SO(B)$ .

$$\Rightarrow T_{A_0} SO(B) = \{A_0 B : B \text{ skew-sym}\}$$

$$= \{A_0 \hat{x} : x \in \mathbb{R}^3\}$$

$$\approx \mathbb{R}^3$$

Since  ~~$A_0^T A_0 = I$~~   $A_0^T \hat{x} = A_0^T \hat{x} A_0$ ,  
 we also have  $T_{A_0} SO(B) = \{A_0 (\hat{A}_0^T x) : x \in \mathbb{R}^3\}$   
 $= \{\hat{x} A_0 : x \in \mathbb{R}^3\}$

Def.  $T_{\mathbb{R}^3} SO(B) \times \mathbb{R}^3 \rightarrow TSO(B)$   
 $T_L(A, x) := A \hat{x}$        $T_R(A, x) := \hat{x} A$

$T_L^{-1}$  is left translation of  $TSO(B)$   
 $T_R^{-1}$  is right translation of  $TSO(B)$

Define ODE on  $SO(B) \times \mathbb{R}^3$ :  
 $\dot{A} = A \underline{\Omega}$

$$\underline{\Pi} \underline{\Omega} = (\underline{\Pi} \underline{\Omega}) \times \underline{\Omega}$$

Alt:  $\dot{A} = A (\underline{\Pi}^{-1} M)$   
 $\dot{M} = M \times (\underline{\Pi}^{-1} M) = (\underline{\Pi}^{-1} M) M$

left view full free rigid body eqs.  
Trivial of 2<sup>nd</sup> order ODE on SO(3).

Important pt: eq for  $\underline{\Omega}$  (or  $\underline{M}$ )  
doesn't depend on  $A$ !

~~Note: Both  $A$  &  $\underline{M}$  have inst angular vel  
 $\underline{\Omega} = \underline{E}^{-2} \underline{M}$ .~~

Note:  $A$  has inst "angular vel"  $\underline{\Omega} = \underline{E}^{-2} \underline{M}$ ;  
 $\underline{M}$  has inst angular vel.  $-\underline{\Omega}$

What happens if you rotate  $\underline{M}$  by  $A$ ?

~~Define  $\underline{\mu} :=$  Given trajectory  $(A(t), \underline{M}(t))$ ,~~

def.  $\underline{\mu}(t) := A(t) \underline{M}(t)$ .

$$\begin{aligned} \dot{\underline{\mu}} &= \dot{A} \underline{M} + A \dot{\underline{M}} \\ &= A \hat{\underline{\Omega}} \underline{M} + A (\underline{M} \times \underline{\Omega}) \\ &= A (\underline{\Omega} \times \underline{M} + \underline{M} \times \underline{\Omega}) = \underline{0}. \end{aligned}$$

Rel ODE:  $\dot{A} = A(\underline{E}^{-2} A^{-2} \underline{a})$ ,

$\underline{\mu}$  const.

If we define  $\rho(A) := A I A^T$ , then  
 $(\rho(A))^{-1} = (A^T)^{-2} \underline{E}^{-2} A^{-2} = A \underline{E}^{-2} A^T$ ;

$$\begin{aligned} \Pi^{-1} A^{-1} \mu &= A^{-1} \underbrace{\mathcal{L}(A)^{-2}}_{\mu} = A^T \mathcal{L}(A)^{-2} \mu \\ \downarrow & \\ A^T \mathcal{L}(A)^{-2} \mu &= A^T \mathcal{L}(A)^{-2} \mu A \\ \Rightarrow \dot{A} &= A (\Pi^{-1} A^{-2} \mu) \\ &= \mathcal{L}(A)^{-2} \mu A \quad \otimes \end{aligned}$$

1st order ODE on  $SO(B)$ , param by  $\mu$ .  
 $\mu$  is spatial ang mom. <sup>-conserved</sup>;  $\mathcal{L}(A)^{-2} \mu$  is spatial ang vel.

Systems sol of one sys map onto sol of other. Neither is fund. supercon.

Note: Conservation of  $\mu$  is directly related to indep of Euler eq from A.

Analogs of the ~~sys~~ trees of systems exist for all Lie groups.  
 Analog of spheres is coadjoint orbits.

What do these eq have to do w rigid bodies?

Consider a 3-D ~~body~~ rigid body — rel. position of pts doesn't change.

Let  $B_{ref}$  denote ~~the region~~ a ref. pos. of the body +  $\varphi_t: B_{ref} \rightarrow \mathbb{R}^3$  denote map to pos at time  $t$ .

Assume  $\varphi_0$  is or. pres.; then  $\varphi_t$  is a Euc motion;  $\exists$  curves  $A(t) \in SO(3)$ ,  $\underline{c}(t) \in \mathbb{R}^3$  s.t.

$$\varphi_t(x) = A(t)x + \underline{c}(t).$$

For convenience, assume ref. coords  $x$  are det. for ref. to center of mass, + that  $\underline{c} \equiv \underline{0}$ . (This involves exp of transl symmetry.)

Let  $\hat{\Omega}_t(t) \in \mathbb{R}^3$  be det. by  $\dot{A}(t) = A(t)\hat{\Omega}_t(t)$ ,

$$\begin{aligned} \text{so } \varphi_t(x) &= A(t)x \\ &= A(t)\hat{\Omega}_t(t)x \\ &= A(t)(\hat{\Omega}_t(t)x). \end{aligned}$$

$$\left\{ \begin{array}{l} KE = \frac{1}{2} \int_{B_{ref}} P(x) \|\dot{\varphi}_t(x)\|^2 dx, \\ \text{where } P \text{ is density.} \end{array} \right.$$

~~Proof~~ Rec,

$$\begin{aligned} \|\varphi(x)\|^2 &= \|A(x)(\underline{Q}(x) \times \underline{x})\|^2 \\ &= \|\underline{Q}(x) \times \underline{x}\|^2 \\ &= \|\underline{Q}(x)\|^2 \|\underline{x}\|^2 - \langle \underline{Q}(x), \underline{x} \rangle^2 \end{aligned}$$

$$\begin{aligned} \Rightarrow E &= \frac{1}{2} \int_{\text{Bref}} P(x) \left( \|\underline{Q}\|^2 \|\underline{x}\|^2 - \langle \underline{Q}, \underline{x} \rangle^2 \right) d\underline{x} \\ &= \frac{1}{2} \langle \underline{Q}, \left( \int_{\text{Bref}} P(x) (\|\underline{x}\|^2 \underline{I} - \underline{x} \underline{x}^T) d\underline{x} \right) \underline{Q} \rangle \end{aligned}$$

Let  $\underline{\Gamma} = \int_{\text{Bref}} P(x) (\|\underline{x}\|^2 \underline{I} - \underline{x} \underline{x}^T) d\underline{x}$ .

(This sets some restrictions on evokes of  $\underline{\Gamma}$ .)