

**Effects of variable viscosity and viscous dissipation on the disappearance of criticality of a reactive third-grade fluid in a slab**

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## **1. The Equations**

**The one-dimensional equations for reactive fully developed flow of an incompressible, thermodynamically compatible fluid of grade three are:**

$$\frac{d}{dy} \left( \exp(-\gamma\theta) \frac{du}{dy} \right) + 6\Lambda \left( \frac{du}{dy} \right)^2 \frac{d^2u}{dy^2} = C, \quad (1)$$

$$\frac{d^2\theta}{dy^2} + \Gamma \left( \frac{du}{dy} \right)^2 \left( \exp(-\gamma\theta) + 2\Lambda \left( \frac{du}{dy} \right)^2 \right) + \delta(1 + \beta\theta)^m \exp\left(\frac{\theta}{1 + \beta\theta}\right) = 0, \quad (2)$$

**together with the corresponding boundary conditions**

$$u(-1) = 0, \quad u(1) = 0, \quad \text{and} \quad \theta(-1) = 0, \quad \theta(1) = 0, \quad (3)$$

**\* We assume heat generation with variable pre-exponential factor (see Massoudi and Phuoc (2008)).**

**\* The viscosity is temperature dependent with  $\gamma \geq 0$  for a liquid and  $\gamma$**

$< 0$  for a gas.

Observation: It is worth noting that the third term in equation (2) represents the heat generation term that has been excluded from Szeri and Rajagopal (1985). In addition, if  $\gamma = 0$ , we are in Okoya 2006.

## 2. Computational method and confirmations

The velocity equation (1) and the temperature equation (2) are coupled and must be solved simultaneously. By means of the shooting method, to some extent modified to cope with given type of non-linearities, the dimensionless governing equations (1) - (2) and the boundary conditions (3) will be numerically investigated with the unknown parameter taken to be  $\delta$ . We are interested in studying the bifurcation of the physical parameters  $(\delta, \beta, \theta_{max})$ , where  $\theta_{max}$  is the value of the maximum temperature of the central plane. It is a well known feature that criticality is represented by discontinuous response of  $\theta_{max}(\delta)$  when  $\beta \neq 0$  (but  $\beta < \beta_{tr}$ ), while criticality disappears for  $\beta = \beta_{tr}$ , with smooth responses of  $\theta_{max}$  to  $\delta$  for higher values of  $\beta$ . A test of the accuracy of the numerical procedure is provided by comparing the results against those for special cases in the literature, i.e. Boddington et al. (1983) and (1984), Shonhiwa and Zaturaska (1987) and Okoya (2006).

## 3. Results

The computed results shown on Figures 1, 2, 3, 4, 5 and 6 (extracted from a forthcoming paper) represent the variation of transitional parameters with respect to the non-Newtonian parameter  $\Lambda$ , viscous heating  $\Gamma$  and the viscosity variational number  $\gamma$

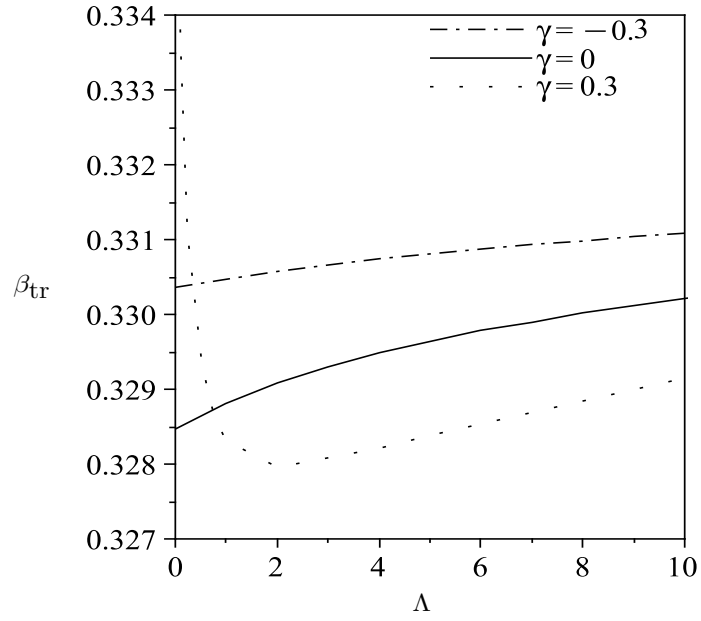


Figure 1: Illustrating the profiles of  $\beta_{tr}$  with  $\Lambda$  for  $C = -0.25$ ,  $m = 0.5$ ,  $\Gamma = 16$ .

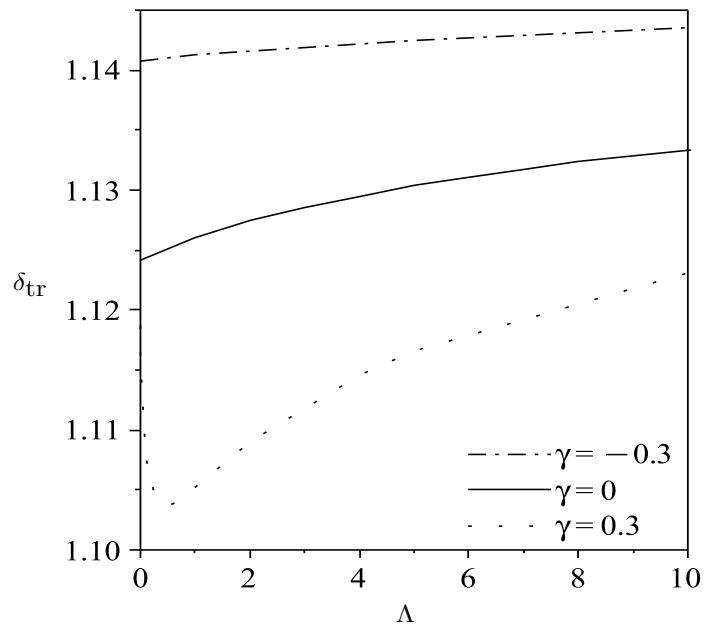


Figure 2: Illustrating the plots of  $\delta_{tr}$  versus  $\Lambda$  for  $C = -0.25$ ,  $m = 0.5$ ,  $\Gamma = 16$ .

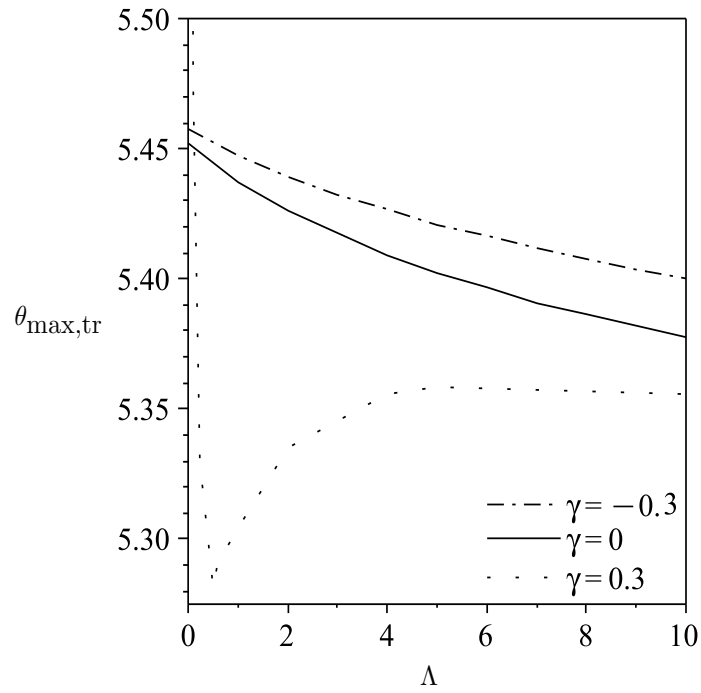


Figure 3: Illustrating the dependence  $\theta_{\max, \text{tr}}$  on  $\Lambda$  for  $C = -0.25$ ,  $m = 0.5$ ,  $\Gamma = 16$ .

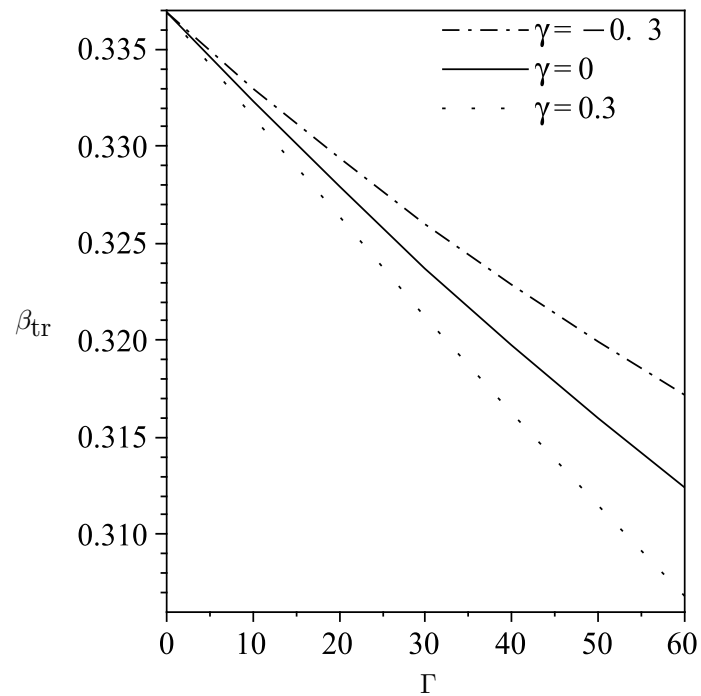


Figure 4: Illustrating the comparisons of  $\beta_{\text{tr}}$  with  $\Gamma$  for  $C = -0.25$ ,  $m = 0.5$  and  $\Lambda = 5$ .

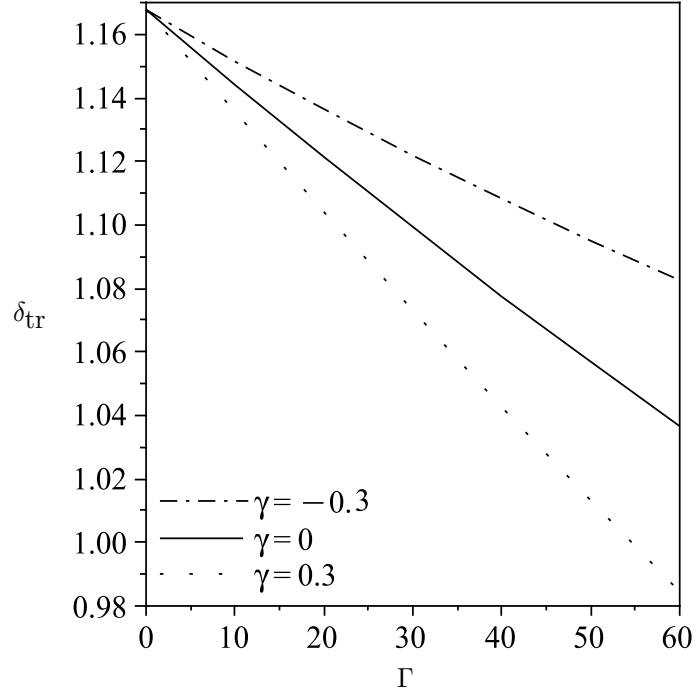


Figure 5: Illustrating the plots of  $\delta_{tr}$  versus  $\Gamma$  for  $C = -0.25$ ,  $m = 0.5$  and  $\Lambda = 5$ .

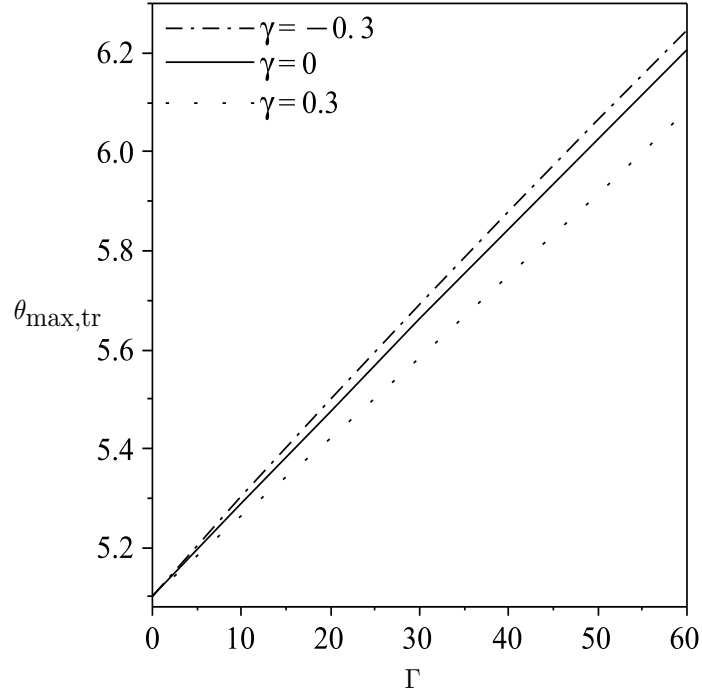


Figure 6: Illustrating the pictures of  $\theta_{max,tr}$  on  $\Gamma$  for  $C = -0.25$ ,  $m = 0.5$  and  $\Lambda = 5$ .

## References

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