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## ADDITIVITY

There are a number of equivalent additivity questions in quantum information theory

We let  $S(\rho) = -\text{Tr} \rho \log \rho$  be the entropy of a positive trace 1 operator

~~Question A~~

Question a)

Is minimum entropy output additive?

Let  $\Phi$  be a quantum channel

$$\Phi(\rho) = \sum_k A_k \rho A_k^\dagger, \text{ with } \sum_k A_k^\dagger A_k = I$$

~~Entropy~~

$$S_{\min}(\Phi) = \min_{|\nu\rangle\langle\nu|} S(\Phi(|\nu\rangle\langle\nu|))$$

$$\text{Is } S_{\min}(\Phi_1) + S_{\min}(\Phi_2) = S_{\min}(\Phi_1 \otimes \Phi_2) \\ \geq \text{easy}$$

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b) Is entanglement of formation additive?

If  $\rho_{AB} \in \mathcal{H}_A \otimes \mathcal{H}_B$  is a density matrix on a tensor product space, ~~and~~ then

$$E_F(\rho_{AB}) = \min_i \sum_i p_i S(\text{Tr}_A |v_i\rangle\langle v_i|)$$

$$\text{where } \rho_{AB} = \sum_i p_i |v_i\rangle\langle v_i|$$

$$\text{Is } E_F(\rho_{AB} \otimes \sigma_{AB}) = E_F(\rho_{AB}) + E_F(\sigma_{AB})$$

$\leq$  is easy

Recall  $S(\text{Tr}_A |v_i\rangle\langle v_i|)$  is the amount of entanglement in a pure state  $|v_i\rangle\langle v_i|$ .

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c) minimum entanglement state in a subspace.

Let  $\Pi$  be a projection matrix in a tensor product space  $\mathcal{H}_A \otimes \mathcal{H}_B$

$$E_{\min}(\Pi) = \min S(\text{Tr}_B(|v\rangle\langle v|))$$
$$|v\rangle\langle v| \leq \Pi$$

Is  $E_{\min}(\Pi_1 \otimes \Pi_2) = E_{\min}(\Pi_1) + E_{\min}(\Pi_2)$   
 $\leq$  easy

d) Strong superadditivity of entanglement of formation.

additivity  ~~$\rho_1, \rho_2 \in \mathcal{H}_A$~~

$$E_F(\rho_1 \otimes \rho_2) = E_F(\rho_1) + E_F(\rho_2)$$

$$\rho_1 \in \mathcal{H}_{A_1} \otimes \mathcal{H}_{B_1}$$

$$\rho_2 \in \mathcal{H}_{A_2} \otimes \mathcal{H}_{B_2}$$

Generalization

$$\text{FOR } \rho \in \mathcal{H}_{A_1} \otimes \mathcal{H}_{A_2} \otimes \mathcal{H}_{B_1} \otimes \mathcal{H}_{B_2}$$

$$\text{Is } E_F(\rho) \geq E_F(\text{Tr}_2 \rho) + E_F(\text{Tr}_1 \rho)$$

For  $E_F$ , trace over  $A_1 \otimes A_2$ ,  $B_1$ , and  $B_2$  respectively.

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So multiplicativity of  $\mathcal{V}_p$

$\Rightarrow$  additivity of min entropy.

Unfortunately, there is a counterexample

$\mathcal{V}_p$  not multiplicative if  $p \geq 5$

In fact, looking at the counterexample for  $p \geq 3$ , there is a local maximum of  $\mathcal{V}_p$  at an entangled input state.

For  $1 \leq p \leq 2$ , everything looks very nice in counterexample

All we need for additivity of entropy is  $\mathcal{V}_p$  multiplicative for  $1 < p \leq 1 + \epsilon$ .

Christopher King

For  $\Delta$  a depolarizing channel

$$\Downarrow \quad \mathcal{V}_p(\Delta \otimes \bar{\Phi}) = \mathcal{V}_p(\Delta) \mathcal{V}_p(\bar{\Phi})$$

( $\bar{\Phi}$  is arbitrary)

$$\Delta(\rho) = \lambda \rho + (1-\lambda) \frac{\mathbb{I}}{d}$$

Also have  $\chi$  additive in some case

⑦

## Entanglement breaking channels

Suppose we have three Hilbert  
space

$\mathcal{H}_{\text{ref}}$

$$\mathcal{H}_{\text{in}} \xrightarrow{\Phi} \mathcal{H}_{\text{out}}$$

$|v\rangle\langle v|$  in  $\mathcal{H}_{\text{in}} \otimes \mathcal{H}_{\text{ref}}$   
possibly entangled

$$\rho = (\Phi \otimes I)(|v\rangle\langle v|) \text{ in } \mathcal{H}_{\text{out}} \otimes \mathcal{H}_{\text{ref}}$$

If  $\rho$  is always  
separable (i.e., not entangled),

then  $\Phi$  is an entanglement breaking  
channel.

For entanglement breaking channels,

$S_{\text{min}}$ ,  $\chi$  are additive  
 $\nu_p$  is multiplicative.

I.e.,  $\Phi$  entanglement breaking,  $\Psi$  arbitrary

$$\nu_p(\Phi \otimes \Psi) = \nu_p(\Phi) \nu_p(\Psi) \quad \chi(\Phi \otimes \Psi) = \chi(\Phi) \cdot \chi(\Psi)$$

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Proof of additivity for e. b. c.,  
for  $S_{\min}$

This uses strong subadditivity of entropy

Take  $\Psi \otimes \bar{\Phi}$ ,  $\Psi$  arbitrary  
 $\bar{\Phi}$  entanglement breaking

$I \otimes \bar{\Phi}(\rho_{AB})$  is separable

$$\text{So } I \otimes \bar{\Phi}(\rho_{AB}) = \sum_j q_j |a_j\rangle\langle a_j| \otimes |b_j\rangle\langle b_j|$$

$$\sigma_{AB} = \Psi \otimes \bar{\Phi}(\rho_{AB}) = \sum_j q_j \Psi(|a_j\rangle\langle a_j|) \otimes |b_j\rangle\langle b_j|$$

Look at state

$$\sigma_{ABC} = \sum_j q_j \Psi(|a_j\rangle\langle a_j|) \otimes |b_j\rangle\langle b_j| \otimes |c_j\rangle\langle c_j|$$

system A                      system B                      system C

We use strong subadditivity on

$\sigma_{ABC}$  to show there is a ~~tensor~~  
~~product~~ output state for  $\bar{\Phi}$  and  $\Psi$   
whose combined entropy at most  
 $S(\sigma_{AB})$

$$(9) \quad S(\sigma_{AB}) \geq S(\sigma_{ABC}) - S(\sigma_{BC}) + S(\sigma_B)$$

$$\begin{aligned} \text{But } \sigma_B &= \sum q_i |b_i\rangle\langle b_i| \\ &= \text{Tr}_A \sum q_i |a_i\rangle\langle a_i| \otimes |b_i\rangle\langle b_i| \\ &= \text{Tr}_A (I \otimes \Phi(\rho_{AB})) \\ &= \Phi(\text{Tr}_A \rho_{AB}) \end{aligned}$$

$$\text{So } S(\sigma_B) \geq S_{\min}(\Phi)$$

Need to show

$$S(\sigma_{ABC}) - S(\sigma_{BC}) \geq S_{\min}(\Psi)$$

$$\sigma_{ABC} = \sum q_j \Psi(|a_j\rangle\langle a_j| \otimes |b_j\rangle\langle b_j| \otimes |j\rangle\langle j|)$$

$$S(\sigma_{ABC}) = S(\sigma_{AC}) \quad \text{since } |b_j\rangle \text{ is uniquely determined by } |j\rangle$$

$$\text{Similarly } S(\sigma_{BC}) = S(\sigma_C)$$

$$\begin{aligned} S(\sigma_{AC}) &= \sum q_i S(\Psi(|a_i\rangle\langle a_i|)) + \sum q_i \log q_i \\ &= S(\sigma_C) \end{aligned}$$

$$\text{So } S(\sigma_{AC}) - S(\sigma_C) \geq S_{\min}(\Psi) \quad \text{QED.}$$

The key ingredient in the proof of multiplicativity for the depolarizing channel is the Lieb-Thirring inequality

$$\text{Tr} \left[ (A^{\frac{1}{2}} B A^{\frac{1}{2}})^p \right] \leq \text{Tr} (A^{p/2} B^p A^{p/2}) = \text{Tr} (A^p B^p)$$

for  $p \geq 1$

Recall the depolarizing channel

$$\Delta_\lambda(\rho) = \lambda \rho + (1-\lambda) I/d$$

We also have the dephasing channel

Suppose  $E_i = |e_i\rangle\langle e_i|$ ,  $|e_i\rangle$  orthonormal basis

Dephasing channel Multiplies off-diag-terms by  $\lambda$ .

$$\Phi_\lambda(\rho) = \lambda \rho + (1-\lambda) \sum_{i=1}^d E_i \rho E_i$$

We can express the depolarizing channel  $\Delta_\lambda$  as a convex combination of dephasing channels

In addition, if our channel is acting on a  $\rho$  with eigenvectors  $|v_j\rangle$ ,

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We can express  $\Delta_\lambda$  as a convex combination of dephasing channels, each of which has the  $E_i$  complementary to  $|v_j\rangle$ .

$$\text{That is, } \langle v_i | E_i | v_j \rangle = \frac{1}{d}$$

This will let us work on the dephasing channel.

Lemma. For a phase damping channel,

$$\|(\Phi_\lambda \otimes I)\rho_{12}\|_p \leq d^{(1-1/p)} \nu_p(\Delta_\lambda) \left[ \sum_{i=1}^d \text{Tr}(\rho_2^{(i)})^p \right]^{1/p}$$

$$\rho_2^{(i)} = \text{Tr}_1 [(E_i \otimes I)\rho_{12}]$$

Lieb Thinning inequality

$$\text{Tr} \left( (A^{1/2} B A^{1/2})^p \right) \leq \text{Tr} (A^p B^p)$$

Clever step: Factor  $(\Phi_\lambda \otimes I)\rho_{12}$  into  $A^{1/2} B A^{1/2}$  so this inequality can be applied.

Use  $\rho_{12} = (I \otimes \Psi)(\sigma_{12})$  to get

~~add~~ multiplicativity

$E_F$  and  $\chi$

Stinespring

Any completely positive trace-preserving  $\Phi$  can be realized

$$\Phi : \rho \rightarrow U(\rho \otimes |0\rangle\langle 0|)U^\dagger$$

$$\downarrow \text{Tr}_2$$

$$(\rho \otimes |0\rangle\langle 0|)U^\dagger$$

Let us call  $U(\rho \otimes |0\rangle\langle 0|)U^\dagger$   
 $\mathcal{U}(\rho)$

$\Phi$  realized as an

- 1) embedding in a larger Hilbert space
- 2) unitary
- 3) partial trace

$\mathcal{U}$  is steps 1 and 2

$$S(\Phi(|\psi\rangle\langle\psi|))$$

$$= S(\text{Tr}_2 \mathcal{U}(|\psi\rangle\langle\psi|))$$

$$= E_F(\mathcal{U}(|\psi\rangle\langle\psi|))$$

$\uparrow$   
pure state entanglement

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entanglement of formation

$$E_F(\mathcal{U}(\rho)) = \min_{\rho = \sum p_i \rho_i} \sum p_i E_F(\mathcal{U}(\rho_i))$$

$$= \min_{\rho = \sum p_i \rho_i} \sum p_i S(\text{Tr}_2 \mathcal{U}(\rho_i))$$

$$= \min_{\rho = \sum p_i \rho_i} \sum p_i S(\Phi(\rho_i))$$

where  $\rho_i$  are pure states.

This is the second term in the Holevo capacity.

This shows  $E_F$  is equivalent to constrained Holevo capacity, where we maximize over states with fixed  $\rho = \sum p_i \rho_i$ .

To go from constrained to unconstrained Holevo capacity, we use technique called Channel Extension



# Equivalence of additivity questions

3 ingredients in proofs

Linear programming  
 Stinespring dilation theorem  
 Channel extension.

Linear programming  
 Need to compute

$$\min_{\rho = \sum_i p_i |v_i\rangle\langle v_i|} \sum p_i S(\Phi(|v_i\rangle\langle v_i|))$$

Assume all  $|v_i\rangle\langle v_i|$ ,  $S(\Phi(|v_i\rangle\langle v_i|))$   
 precomputed in big table

$$\min_{p_i \geq 0} \sum p_i S(\Phi(|v_i\rangle\langle v_i|))$$

subject to

$$\sum_i p_i |v_i\rangle\langle v_i| = \rho$$

Dual LP

$$\max_{\tau} \text{Tr } \tau \rho$$

subject to

$$\langle v | \tau | v \rangle \leq S(\Phi(|v\rangle\langle v|)) \quad \forall |v\rangle \in \mathcal{A} \otimes \mathcal{B}$$

$\tau$  is maximized over Hermitian matrices.

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Dual LP has same optimum  
as primal

Proof one direction is easy

$$\begin{aligned} \min \sum_i p_i S(\Phi(|v_i\rangle\langle v_i|)) &\geq \sum_i p_i \langle v_i | \tau | v_i \rangle \\ &= \text{Tr} \sum_i p_i \tau |v_i\rangle\langle v_i| \\ &= \text{Tr} \tau \rho \end{aligned}$$

Other direction needs to  
involve duality theorem.

Additivity of min entropy  $\Rightarrow$   
additivity of ~~of~~ constrained  
Holevo capacity (fix  $\chi$ )

$$\chi_\rho(\Phi) = \max_{\sum p_i \rho_i = \rho} S(\Phi(\rho)) - \sum_i p_i S(\Phi(\rho_i))$$

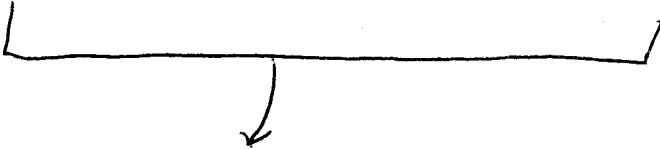
By LP, there is  $\tau$  so.

$$\text{Tr} \rho_i \tau \leq S(\Phi(\rho_i))$$

with equality for signal states  $\rho_i$

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Find another channel  $\tilde{\Phi}$

$$s.o. \quad S(\tilde{\Phi}(\rho)) = S(\Phi(\rho)) + C - \text{Tr} \rho \alpha$$


This side  $\geq C$  for  
signal states  $\rho_i$

$\geq C$  for other states

So signal states for this channel  
are also minimum entropy output states

If we have additivity of min entropy  
output for this channel,  
gives additivity of  $X_{(\text{constrained})}$  for this  
channel  
gives additivity of  $X(\text{constrained})$   
for the original channel.

Cannot find such a channel.

Can find a sequence of ~~such~~ channels  
so it works "in the limit".

(17) Additivity of  $\chi$  also proved for a number of other channels

unital qubit channels

Unital means  $\Phi(I/2) = I/2$ .